

$$I^G(J^P) = 1^-(0^-)$$

We have omitted some results that have been superseded by later experiments. The omitted results may be found in our 1988 edition Physics Letters **B204** 1 (1988).

π^\pm MASS

The most accurate charged pion mass measurements are based upon x-ray wavelength measurements for transitions in π^- -mesonic atoms. The observed line is the blend of three components, corresponding to different K-shell occupancies. JECKELMANN 94 revisits the occupancy question, with the conclusion that two sets of occupancy ratios, resulting in two different pion masses (Solutions A and B), are equally probable. We choose the higher Solution B since only this solution is consistent with a positive mass-squared for the muon neutrino, given the precise muon momentum measurements now available (DAUM 91, ASSAMAGAN 94, and ASSAMAGAN 96) for the decay of pions at rest. Earlier mass determinations with pi-mesonic atoms may have used incorrect K-shell screening corrections.

Measurements with an error of > 0.005 MeV have been omitted from this Listing.

| <u>VALUE (MeV)</u> | <u>DOCUMENT ID</u> | <u>TECN</u> | <u>CHG</u> | <u>COMMENT</u> |
|---|---|-------------|------------|--|
| 139.57061 ± 0.00024 OUR FIT | Error includes scale factor of 1.6. | | | |
| 139.57061 ± 0.00023 OUR AVERAGE | Error includes scale factor of 1.5. See the ideogram below. | | | |
| 139.57077 ± 0.00018 | ¹ TRASSINELLI 16 | CNTR | | X-ray transitions in pionic N ₂ |
| 139.57071 ± 0.00053 | ² LENZ 98 | CNTR | – | pionic N ₂ -atoms gas target |
| 139.56995 ± 0.00035 | ³ JECKELMANN 94 | CNTR | – | π^- atom, Soln. B |
| ● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ● | | | | |
| 139.57022 ± 0.00014 | ⁴ ASSAMAGAN 96 | SPEC | + | $\pi^+ \rightarrow \mu^+ \nu_\mu$ |
| 139.56782 ± 0.00037 | ⁵ JECKELMANN 94 | CNTR | – | π^- atom, Soln. A |
| 139.56996 ± 0.00067 | ⁶ DAUM 91 | SPEC | + | $\pi^+ \rightarrow \mu^+ \nu$ |
| 139.56752 ± 0.00037 | ⁷ JECKELMANN 86B | CNTR | – | Mesonic atoms |
| 139.5704 ± 0.0011 | ⁶ ABELA 84 | SPEC | + | See DAUM 91 |
| 139.5664 ± 0.0009 | ⁸ LU 80 | CNTR | – | Mesonic atoms |
| 139.5686 ± 0.0020 | CARTER 76 | CNTR | – | Mesonic atoms |
| 139.5660 ± 0.0024 | ^{8,9} MARUSHEN... 76 | CNTR | – | Mesonic atoms |

¹ TRASSINELLI 16 use the muonic oxygen line for online energy calibration of the pionic line.

² LENZ 98 result does not suffer K-electron configuration uncertainties as does JECKELMANN 94.

³ JECKELMANN 94 Solution B (dominant 2-electron K-shell occupancy), chosen for consistency with positive $m_{\nu_\mu}^2$.

⁴ ASSAMAGAN 96 measures the μ^+ momentum p_μ in $\pi^+ \rightarrow \mu^+ \nu_\mu$ decay at rest to be 29.79200 ± 0.00011 MeV/c. Combined with the μ^+ mass and the assumption $m_{\nu_\mu} = 0$, this gives the π^+ mass above; if $m_{\nu_\mu} > 0$, m_{π^+} given above is a lower limit.

Combined instead with m_μ and (assuming *CPT*) the π^- mass of JECKELMANN 94, p_μ gives an upper limit on m_{ν_μ} (see the ν_μ).

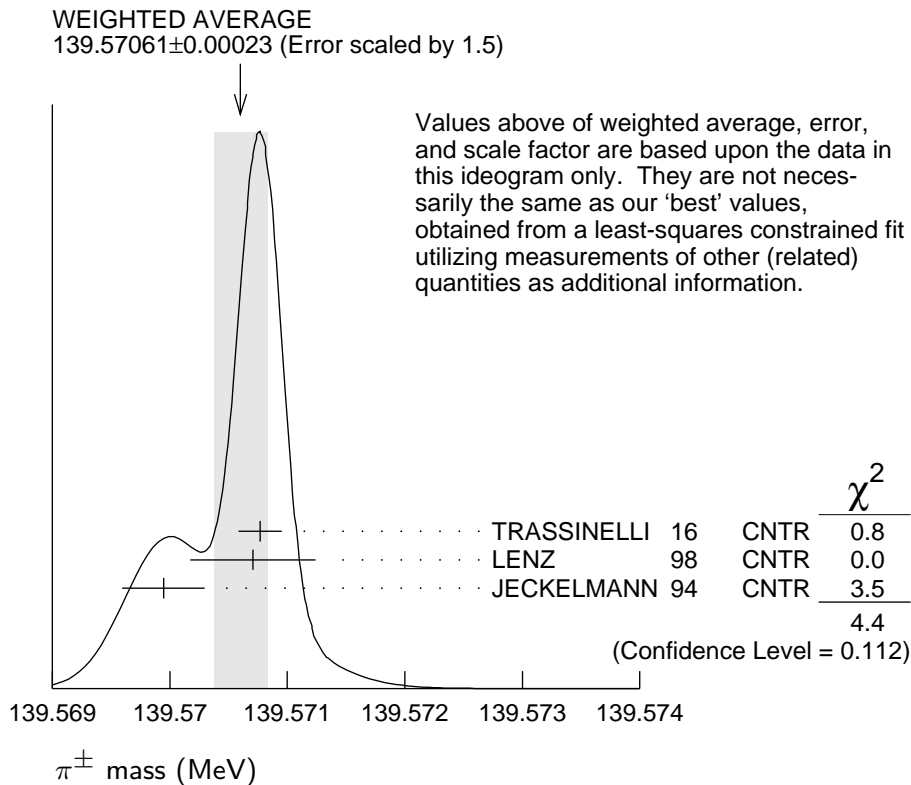
⁵ JECKELMANN 94 Solution A (small 2-electron K-shell occupancy) in combination with either the DAUM 91 or ASSAMAGAN 94 pion decay muon momentum measurement yields a significantly negative $m_{\nu_\mu}^2$. It is accordingly not used in our fits.

⁶ The DAUM 91 value includes the ABELA 84 result. The value is based on a measurement of the μ^+ momentum for π^+ decay at rest, $p_\mu = 29.79179 \pm 0.00053$ MeV, uses $m_\mu = 105.658389 \pm 0.000034$ MeV, and assumes that $m_{\nu_\mu} = 0$. The last assumption means that in fact the value is a lower limit.

⁷ JECKELMANN 86B gives $m_\pi/m_e = 273.12677(71)$. We use $m_e = 0.51099906(15)$ MeV from COHEN 87. The authors note that two solutions for the probability distribution of K-shell occupancy fit equally well, and use other data to choose the lower of the two possible π^\pm masses.

⁸ These values are scaled with a new wavelength-energy conversion factor $V\lambda = 1.23984244(37) \times 10^{-6}$ eV m from COHEN 87. The LU 80 screening correction relies upon a theoretical calculation of inner-shell refilling rates.

⁹ This MARUSHENKO 76 value used at the authors' request to use the accepted set of calibration γ energies. Error increased from 0.0017 MeV to include QED calculation error of 0.0017 MeV (12 ppm).



$m_{\pi^+} - m_{\mu^+}$

Measurements with an error > 0.05 MeV have been omitted from this Listing.

VALUE (MeV) EVTS DOCUMENT ID TECN CHG COMMENT

● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●

| | | | | | | |
|--------------------|-----|-------------------|----|------|---|-------------------------------|
| 33.91157 ± 0.00067 | | ¹ DAUM | 91 | SPEC | + | $\pi^+ \rightarrow \mu^+ \nu$ |
| 33.9111 ± 0.0011 | | ABELA | 84 | SPEC | | See DAUM 91 |
| 33.925 ± 0.025 | | BOOTH | 70 | CNTR | + | Magnetic spect. |
| 33.881 ± 0.035 | 145 | HYMAN | 67 | HEBC | + | K^- He |

¹ The DAUM 91 value assumes that $m_{\nu_\mu} = 0$ and uses our $m_\mu = 105.658389 \pm 0.000034$ MeV.

$$(m_{\pi^+} - m_{\pi^-}) / m_{\text{average}}$$

A test of *CPT* invariance.

| <u>VALUE (units 10⁻⁴)</u> | <u>DOCUMENT ID</u> | <u>TECN</u> |
|--------------------------------------|--------------------|-------------|
| 2 ± 5 | AYRES | 71 CNTR |

π^\pm MEAN LIFE

Measurements with an error > 0.02×10^{-8} s have been omitted.

| <u>VALUE (10⁻⁸ s)</u> | <u>DOCUMENT ID</u> | <u>TECN</u> | <u>CHG</u> | <u>COMMENT</u> | |
|---|---------------------|-------------------------------------|------------|----------------|--------------------|
| 2.6033 ± 0.0005 | OUR AVERAGE | Error includes scale factor of 1.2. | | | |
| 2.60361 ± 0.00052 | ¹ KOPTEV | 95 | SPEC | + | Surface μ^+ 's |
| 2.60231 ± 0.00050 ± 0.00084 | NUMAO | 95 | SPEC | + | Surface μ^+ 's |
| 2.609 ± 0.008 | DUNAITSEV | 73 | CNTR | + | |
| 2.602 ± 0.004 | AYRES | 71 | CNTR | ± | |
| 2.604 ± 0.005 | NORDBERG | 67 | CNTR | + | |
| 2.602 ± 0.004 | ECKHAUSE | 65 | CNTR | + | |
| ● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ● | | | | | |
| 2.640 ± 0.008 | ² KINSEY | 66 | CNTR | + | |

¹ KOPTEV 95 combines the statistical and systematic errors; the statistical error dominates.

² Systematic errors in the calibration of this experiment are discussed by NORDBERG 67.

$$(\tau_{\pi^+} - \tau_{\pi^-}) / \tau_{\text{average}}$$

A test of *CPT* invariance.

| <u>VALUE (units 10⁻⁴)</u> | <u>DOCUMENT ID</u> | <u>TECN</u> |
|---|------------------------|-------------|
| 5.5 ± 7.1 | AYRES | 71 CNTR |
| ● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ● | | |
| -14 ± 29 | PETRUKHIN | 68 CNTR |
| 40 ± 70 | BARDON | 66 CNTR |
| 23 ± 40 | ¹ LOBKOWICZ | 66 CNTR |

¹ This is the most conservative value given by LOBKOWICZ 66.

π ELECTRIC POLARIZABILITY α_π

See HOLSTEIN 14 for a general review on hadron polarizability.

| <u>VALUE (10⁻⁴ fm³)</u> | <u>EVTS</u> | <u>DOCUMENT ID</u> | <u>TECN</u> | <u>COMMENT</u> |
|---|-------------|---------------------|-------------|---|
| 2.0 ± 0.6 ± 0.7 | 63k | ¹ ADOLPH | 15A | SPEC $\pi^- \gamma \rightarrow \pi^- \gamma$ Compton scatt. |

¹Value is derived assuming $\alpha_\pi = -\beta_\pi$.

π^+ DECAY MODES

π^- modes are charge conjugates of the modes below.

For decay limits to particles which are not established, see the section on Searches for Axions and Other Very Light Bosons.

| Mode | Fraction (Γ_i/Γ) | Confidence level |
|--------------------------------------|--|------------------|
| Γ_1 $\mu^+ \nu_\mu$ | [a] $(99.98770 \pm 0.00004) \%$ | |
| Γ_2 $\mu^+ \nu_\mu \gamma$ | [b] $(2.00 \pm 0.25) \times 10^{-4}$ | |
| Γ_3 $e^+ \nu_e$ | [a] $(1.230 \pm 0.004) \times 10^{-4}$ | |
| Γ_4 $e^+ \nu_e \gamma$ | [b] $(7.39 \pm 0.05) \times 10^{-7}$ | |
| Γ_5 $e^+ \nu_e \pi^0$ | $(1.036 \pm 0.006) \times 10^{-8}$ | |
| Γ_6 $e^+ \nu_e e^+ e^-$ | $(3.2 \pm 0.5) \times 10^{-9}$ | |
| Γ_7 $e^+ \nu_e \nu \bar{\nu}$ | $< 5 \times 10^{-6}$ | 90% |

Lepton Family number (LF) or Lepton number (L) violating modes

| | | | | |
|-----------------------------------|------|-------------|------------------|-----|
| Γ_8 $\mu^+ \bar{\nu}_e$ | L | [c] < 1.5 | $\times 10^{-3}$ | 90% |
| Γ_9 $\mu^+ \nu_e$ | LF | [c] < 8.0 | $\times 10^{-3}$ | 90% |
| Γ_{10} $\mu^- e^+ e^+ \nu$ | LF | < 1.6 | $\times 10^{-6}$ | 90% |

[a] Measurements of $\Gamma(e^+ \nu_e)/\Gamma(\mu^+ \nu_\mu)$ always include decays with γ 's, and measurements of $\Gamma(e^+ \nu_e \gamma)$ and $\Gamma(\mu^+ \nu_\mu \gamma)$ never include low-energy γ 's. Therefore, since no clean separation is possible, we consider the modes with γ 's to be subreactions of the modes without them, and let $[\Gamma(e^+ \nu_e) + \Gamma(\mu^+ \nu_\mu)]/\Gamma_{\text{total}} = 100\%$.

[b] See the Particle Listings below for the energy limits used in this measurement; low-energy γ 's are not included.

[c] Derived from an analysis of neutrino-oscillation experiments.

π^+ BRANCHING RATIOS

$\Gamma(e^+ \nu_e)/\Gamma_{\text{total}}$

Γ_3/Γ

See note [a] in the list of π^+ decay modes just above, and see also the next block of data. See also the note on "Decay Constants of Charged Pseudoscalar Mesons" in the D_s^+ Listings.

VALUE (units 10^{-4})

DOCUMENT ID

1.230 \pm 0.004 OUR EVALUATION

$$\frac{[\Gamma(e^+ \nu_e) + \Gamma(e^+ \nu_e \gamma)]}{[\Gamma(\mu^+ \nu_\mu) + \Gamma(\mu^+ \nu_\mu \gamma)]} \quad (\Gamma_3 + \Gamma_4) / (\Gamma_1 + \Gamma_2)$$

See note [a] in the list of π^+ decay modes above. See NUMAO 92 for a discussion of e - μ universality. See also the note on "Decay Constants of Charged Pseudoscalar Mesons" in the D_s^+ Listings.

| VALUE (units 10^{-4}) | EVTS | DOCUMENT ID | TECN | CHG | COMMENT |
|---|------|-------------------------|------|-----|------------------|
| 1.2327 ± 0.0023 OUR AVERAGE | | | | | |
| 1.2344 ± 0.0023 ± 0.0019 | 400k | AGUILAR-AR...15 | CNTR | + | Stopping π^+ |
| 1.2346 ± 0.0035 ± 0.0036 | 120k | CZAPEK 93 | CALO | | Stopping π^+ |
| 1.2265 ± 0.0034 ± 0.0044 | 190k | BRITTON 92 | CNTR | | Stopping π^+ |
| 1.218 ± 0.014 | 32k | BRYMAN 86 | CNTR | | Stopping π^+ |
| • • • We do not use the following data for averages, fits, limits, etc. • • • | | | | | |
| 1.273 ± 0.028 | 11k | ¹ DICAPUA 64 | CNTR | | |
| 1.21 ± 0.07 | | ANDERSON 60 | SPEC | | |

¹ DICAPUA 64 has been updated using the current mean life.

$$\frac{\Gamma(\mu^+ \nu_\mu \gamma)}{\Gamma_{\text{total}}} \quad \Gamma_2 / \Gamma$$

Note that measurements here do not cover the full kinematic range.

| VALUE (units 10^{-4}) | EVTS | DOCUMENT ID | TECN | CHG | COMMENT |
|---|------|------------------------|------|-----|-----------------------------|
| 2.0 ± 0.24 ± 0.08 | | | | | |
| | | ¹ BRESSI 98 | CALO | + | Stopping π^+ |
| • • • We do not use the following data for averages, fits, limits, etc. • • • | | | | | |
| 1.24 ± 0.25 | 26 | CASTAGNOLI 58 | EMUL | | $KE_\mu < 3.38 \text{ MeV}$ |

¹ BRESSI 98 result is given for $E_\gamma > 1 \text{ MeV}$ only. Result agrees with QED expectation, 2.283×10^{-4} and does not confirm discrepancy of earlier experiment CASTAGNOLI 58.

$$\frac{\Gamma(e^+ \nu_e \gamma)}{\Gamma_{\text{total}}} \quad \Gamma_4 / \Gamma$$

The very different values reflect the very different kinematic ranges covered (bigger range, bigger value). And none of them covers the whole kinematic range.

| VALUE (units 10^{-8}) | EVTS | DOCUMENT ID | TECN | CHG | COMMENT |
|---|------|--------------------------|------|-----|---|
| 73.86 ± 0.54 | | | | | |
| | 65k | ¹ BYCHKOV 09 | PIBE | | $e^+ \nu_\gamma$ at rest |
| • • • We do not use the following data for averages, fits, limits, etc. • • • | | | | | |
| 16.1 ± 2.3 | | ² BOLOTOV 90B | SPEC | | 17 GeV $\pi^- \rightarrow e^- \bar{\nu}_e \gamma$ |
| 5.6 ± 0.7 | 226 | ³ STETZ 78 | SPEC | | $P_e > 56 \text{ MeV}/c$ |
| 3.0 | 143 | DEPOMMIER 63B | CNTR | | $(KE)_{e+\gamma} > 48 \text{ MeV}$ |

¹ This BYCHKOV 09 value is for $E_\gamma > 10 \text{ MeV}$ and $\Theta_{e+\gamma} > 40^\circ$.

² BOLOTOV 90B is for $E_\gamma > 21 \text{ MeV}$, $E_e > 70 - 0.8 E_\gamma$.

³ STETZ 78 is for an $e^- \gamma$ opening angle $> 132^\circ$. Obtains 3.7 when using same cutoffs as DEPOMMIER 63B.

$$\frac{\Gamma(e^+ \nu_e \pi^0)}{\Gamma_{\text{total}}} \quad \Gamma_5 / \Gamma$$

| VALUE (units 10^{-8}) | EVTS | DOCUMENT ID | TECN | CHG | COMMENT |
|----------------------------------|------|---------------------------|------|-----|---------------------|
| 1.036 ± 0.006 OUR AVERAGE | | | | | |
| 1.036 ± 0.006 | 64k | ^{1,2} POCANIC 04 | PIBE | + | π decay at rest |
| 1.026 ± 0.039 | 1224 | ³ MCFARLANE 85 | CNTR | + | Decay in flight |
| 1.00 +0.08 -0.10 | 332 | DEPOMMIER 68 | CNTR | + | |
| 1.07 ± 0.21 | 38 | ⁴ BACASTOW 65 | OSPK | + | |

| | | | | | |
|------------|----|------------------------|----|------|---|
| 1.10 ±0.26 | | ⁴ BERTRAM | 65 | OSPK | + |
| 1.1 ±0.2 | 43 | ⁴ DUNAITSEV | 65 | CNTR | + |
| 0.97 ±0.20 | 36 | ⁴ BARTLETT | 64 | OSPK | + |

• • • We do not use the following data for averages, fits, limits, etc. • • •

| | | | | | | |
|------------|----|------------------------|----|------|---|------------------|
| 1.15 ±0.22 | 52 | ⁴ DEPOMMIER | 63 | CNTR | + | See DEPOMMIER 68 |
|------------|----|------------------------|----|------|---|------------------|

¹ POCANIC 04 normalizes to $e^+ \nu_e$ decays, using the PDG 2004 value $B(\pi^+ \rightarrow e^+ \nu_e) = (1.230 \pm 0.004) \times 10^{-4}$. We add their statistical (0.004×10^{-8}), systematic (0.004×10^{-8}) and systematic error due to the uncertainty of $B(\pi^+ \rightarrow e^+ \nu_e)$ (0.003×10^{-8}) in quadrature.

² This result can be used to calculate V_{ud} from pion beta decay: $V_{ud}^{PIBETA} = 0.9728 \pm 0.0030$.

³ MCFARLANE 85 combines a measured rate (0.394 ± 0.015)/s with 1982 PDG mean life.

⁴ DEPOMMIER 68 says the result of DEPOMMIER 63 is at least 10% too large because of a systematic error in the π^0 detection efficiency, and that this may be true of all the previous measurements (also V. Soergel, private communication, 1972).

$\Gamma(e^+ \nu_e e^+ e^-) / \Gamma(\mu^+ \nu_\mu)$ Γ_6 / Γ_1

| VALUE (units 10^{-9}) | CL% | EVTS | DOCUMENT ID | TECN | COMMENT |
|--------------------------|-----|------|-------------|------|--|
| 3.2 ±0.5 ±0.2 | | 98 | EGLI | 89 | SPEC Uses $R_{PCAC} = 0.068 \pm 0.004$ |

• • • We do not use the following data for averages, fits, limits, etc. • • •

| | | | | | |
|------------------|----|----------------------|-----|------|-----------------|
| 0.46 ±0.16 ±0.07 | 7 | ¹ BARANOV | 92 | SPEC | Stopped π^+ |
| < 4.8 | 90 | KORENCHE... | 76B | SPEC | |
| <34 | 90 | KORENCHE... | 71 | OSPK | |

¹ This measurement by BARANOV 92 is of the structure-dependent part of the decay. The value depends on values assumed for ratios of form factors.

$\Gamma(e^+ \nu_e \nu \bar{\nu}) / \Gamma_{total}$ Γ_7 / Γ

| VALUE (units 10^{-6}) | CL% | DOCUMENT ID | TECN | |
|--------------------------|-----|-------------|------|------|
| <5 | 90 | PICCIOTTO | 88 | SPEC |

$\Gamma(\mu^+ \bar{\nu}_e) / \Gamma_{total}$ Γ_8 / Γ

Forbidden by total lepton number conservation. See the note on “Decay Constants of Charged Pseudoscalar Mesons” in the D_s^+ Listings.

| VALUE (units 10^{-3}) | CL% | DOCUMENT ID | TECN | COMMENT |
|--------------------------|-----|---------------------|------|--------------------------|
| <1.5 | 90 | ¹ COOPER | 82 | HLBC Wideband ν beam |

¹ COOPER 82 limit on $\bar{\nu}_e$ observation is here interpreted as a limit on lepton number violation.

$\Gamma(\mu^+ \nu_e) / \Gamma_{total}$ Γ_9 / Γ

Forbidden by lepton family number conservation.

| VALUE (units 10^{-3}) | CL% | DOCUMENT ID | TECN | COMMENT |
|--------------------------|-----|---------------------|------|--------------------------|
| <8.0 | 90 | ¹ COOPER | 82 | HLBC Wideband ν beam |

¹ COOPER 82 limit on ν_e observation is here interpreted as a limit on lepton family number violation.

$\Gamma(\mu^- e^+ e^+ \nu)/\Gamma_{\text{total}}$ Γ_{10}/Γ

Forbidden by lepton family number conservation.

| VALUE (units 10^{-6}) | CL% | DOCUMENT ID | TECN | CHG |
|---|-----|-------------|------|--------|
| <1.6 | 90 | BARANOV | 91B | SPEC + |
| ● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ● | | | | |
| <7.7 | 90 | KORENCHE... | 87 | SPEC + |

 π^+ — POLARIZATION OF EMITTED μ^+ $\pi^+ \rightarrow \mu^+ \nu$

Tests the Lorentz structure of leptonic charged weak interactions.

| VALUE | CL% | DOCUMENT ID | TECN | CHG | COMMENT |
|---|-----|-----------------------|------|------|----------------|
| ● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ● | | | | | |
| <(-0.9959) | 90 | ¹ FETSCHER | 84 | RVUE | + |
| -0.99 ± 0.16 | | ² ABELA | 83 | SPEC | - μ X-rays |

¹FETSCHER 84 uses only the measurement of CARR 83.²Sign of measurement reversed in ABELA 83 to compare with μ^+ measurements.

See the related review(s):

[Form Factors for Radiative Pion and Kaon Decays](#) π^\pm FORM FACTORS F_V , VECTOR FORM FACTOR

| VALUE | EVTS | DOCUMENT ID | TECN | COMMENT |
|--|------|----------------------|------|--|
| 0.0254 ± 0.0017 OUR AVERAGE | | | | |
| 0.0258 ± 0.0017 | 65k | ¹ BYCHKOV | 09 | PIBE $e^+ \nu \gamma$ at rest |
| 0.014 ± 0.009 | | ² BOLOTOV | 90B | SPEC 17 GeV $\pi^- \rightarrow e^- \bar{\nu}_e \gamma$ |
| 0.023 $\begin{smallmatrix} +0.015 \\ -0.013 \end{smallmatrix}$ | 98 | EGLI | 89 | SPEC $\pi^+ \rightarrow e^+ \nu_e e^+ e^-$ |

¹The BYCHKOV 09 F_A and F_V results are highly (anti-)correlated: $F_A + 1.0286 F_V = 0.03853 \pm 0.00014$.²BOLOTOV 90B only determines the absolute value. F_A , AXIAL-VECTOR FORM FACTOR

| VALUE | EVTS | DOCUMENT ID | TECN | COMMENT |
|---|------|-------------------------|------|--|
| 0.0119 ± 0.0001 | | | | |
| ● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ● | | | | |
| 0.0115 ± 0.0004 | 41k | ^{1,3} FRLEZ | 04 | PIBE $\pi^+ \rightarrow e^+ \nu \gamma$ at rest |
| 0.0106 ± 0.0060 | | ^{1,4} BOLOTOV | 90B | SPEC 17 GeV $\pi^- \rightarrow e^- \bar{\nu}_e \gamma$ |
| 0.021 $\begin{smallmatrix} +0.011 \\ -0.013 \end{smallmatrix}$ | 98 | EGLI | 89 | SPEC $\pi^+ \rightarrow e^+ \nu_e e^+ e^-$ |
| 0.0135 ± 0.0016 | | ^{1,4} BAY | 86 | SPEC $\pi^+ \rightarrow e^+ \nu \gamma$ |
| 0.006 ± 0.003 | | ^{1,4} PIILONEN | 86 | SPEC $\pi^+ \rightarrow e^+ \nu \gamma$ |
| 0.011 ± 0.003 | | ^{1,4,5} STETZ | 78 | SPEC $\pi^+ \rightarrow e^+ \nu \gamma$ |

- ¹ These values come from fixing the vector form factor at the CVC prediction, $F_V = 0.0259 \pm 0.0005$.
² When F_V is released, the BYCHKOV 09 F_A is 0.0117 ± 0.0017 , and F_A and F_V results are highly (anti-)correlated: $F_A + 1.0286 F_V = 0.03853 \pm 0.00014$.
³ The sign of $\gamma = F_A / F_V$ is determined to be positive.
⁴ Only the absolute value of F_A is determined.
⁵ The result of STETZ 78 has a two-fold ambiguity. We take the solution compatible with later determinations.

VECTOR FORM FACTOR SLOPE PARAMETER a

This is a in $F_V(q^2) = F_V(0) (1 + a q^2)$

| VALUE | EVTS | DOCUMENT ID | TECN | COMMENT |
|--------------------|------|-------------|------|--------------------------|
| 0.10 ± 0.06 | 65k | BYCHKOV 09 | PIBE | $e^+ \nu \gamma$ at rest |

R, SECOND AXIAL-VECTOR FORM FACTOR

| VALUE | EVTS | DOCUMENT ID | TECN | COMMENT |
|--|------|-------------|------|---------------------------------------|
| 0.059^{+0.009}_{-0.008} | 98 | EGLI 89 | SPEC | $\pi^+ \rightarrow e^+ \nu_e e^+ e^-$ |

π^\pm CHARGE RADIUS

The charge radius of the pion $\sqrt{\langle r_\pi^2 \rangle}$ is defined in relation to the form factor of the pion electromagnetic vertex, called vector form factor VFF, F_π^V . The VFF is a function of the squared four-momentum transfer t , or of the squared c.m. energy s , depending on the channel in which the photon exchange takes place. In both cases, it is related to the slope of the VFF at zero, namely

$$\langle r_\pi^2 \rangle = 6 \frac{dF_\pi^V(q)}{dq} (q=0) \text{ where } q = t, s.$$

The quantity cannot be measured directly. It can be extracted from the cross sections of three processes: pion electroproduction, $eN \rightarrow eN\pi$, and pion electron scattering $e\pi \rightarrow e\pi$, for the t channel, and positron electron annihilation into two charged pions, $e^+e^- \rightarrow \pi^+\pi^-$, for the s channel. We encode all measurements, but we do not use electroproduction data in averaging because the extraction of the pion radius involves, in this case, theoretical uncertainties that cannot be controlled at the needed level of accuracy. In case of analyses based on the same data set, as ANANTHANARAYAN 17 and COLANGELO 19, which cannot be averaged, we combine the results into a common value, with the uncertainty range chosen to cover both analyses. Note that for consistency the form factor needs to be defined in both channels with the vacuum polarisation removed. For details see COLANGELO 19 or Appendix B of ANANTHANARAYAN 16A.

| VALUE (fm) | DOCUMENT ID | TECN | COMMENT |
|----------------------------------|---------------------|------|---------------------------|
| 0.659 ± 0.004 OUR AVERAGE | | | |
| 0.656 ± 0.005 | ¹ PDG 19 | FIT | |
| 0.65 ± 0.05 ± 0.06 | ESCHRICH 01 | CNTR | $\pi e \rightarrow \pi e$ |
| 0.663 ± 0.006 | AMENDOLIA 86 | CNTR | $\pi e \rightarrow \pi e$ |
| 0.663 ± 0.023 | DALLY 82 | CNTR | $\pi e \rightarrow \pi e$ |

• • • We do not use the following data for averages, fits, limits, etc. • • •

| | | | |
|---|-----------------------------|------|-----------------------------------|
| 0.655 ± 0.004 | ² COLANGELO 19 | FIT | Fit existing data |
| 0.657 ± 0.003 | ³ ANANTHANA...17 | FIT | Fit existing data |
| 0.6603 ± 0.0005 ± 0.0004 | ⁴ HANHART 17 | FIT | Fit existing data |
| 0.740 ± 0.031 | ⁵ LIESENFELD 99 | CNTR | $e p \rightarrow e \pi^+ n$ |
| 0.661 ± 0.012 | ⁶ BIJNENS 98 | CNTR | χ PT extraction |
| 0.660 ± 0.024 | AMENDOLIA 84 | CNTR | $\pi e \rightarrow \pi e$ |
| 0.711 ± 0.009 ± 0.016 | ⁵ BEBEK 78 | CNTR | $e N \rightarrow e \pi N$ |
| 0.678 ± 0.004 ± 0.008 | ⁷ QUENZER 78 | CNTR | $e^+ e^- \rightarrow \pi^+ \pi^-$ |
| 0.78 ^{+0.09} _{-0.10} | ADYLOV 77 | CNTR | $\pi e \rightarrow \pi e$ |
| 0.74 ^{+0.11} _{-0.13} | BARDIN 77 | CNTR | $e p \rightarrow e \pi^+ n$ |
| 0.56 ± 0.04 | DALLY 77 | CNTR | $\pi e \rightarrow \pi e$ |

¹ This value combines the measurements of ANANTHANARAYAN 17 and COLANGELO 19 which are based on the same data set. The uncertainty range is chosen to cover both results.

² COLANGELO 19 fit existing F_V data, using an extended Omnes dispersive representation. This analysis is based on the same data set of ANANTHANARAYAN 17. Accordingly, they cannot be averaged. We combine the results into a common value, with the uncertainty range chosen to cover the uncertainty ranges of both analyses.

³ ANANTHANARAYAN 17 fit existing F_V data, using a mixed phase-modulus dispersive representation. This analysis is based on the same data set of COLANGELO 19. Accordingly, they cannot be averaged. We combine the results into a common value, with the uncertainty range chosen to cover the uncertainty ranges of both analyses.

⁴ According to the authors the uncertainty could be underestimated. The value quoted omits the BaBar data AUBERT 09.

⁵ The extractions could contain an additional theoretical uncertainty which cannot be sufficiently quantified.

⁶ BIJNENS 98 fits existing data.

⁷ The extraction is based on a parametrization that does not have correct analytic properties.

π^\pm REFERENCES

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