## $b^{\prime}$ (4 ${ }^{\text {th }}$ Generation) Quark, Searches for

## $b^{\prime}(-1 / 3)$-quark/hadron mass limits in $p \bar{p}$ and $p p$ collisions

| VALUE (GeV) | CL\% | DOCUMENT ID | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: |
| >1570 | 95 | 1 SIRUNYAN 20 | 20BI CMS | $\mathrm{B}\left(b^{\prime} \rightarrow H b\right)=1$ |
| $>1390$ | 95 | 1 SIRUNYAN 2 | 20BI CMS | $\mathrm{B}\left(b^{\prime} \rightarrow Z b\right)=1$ |
| $>1130$ | 95 | 2 SIRUNYAN | 19AQ CMS | $\mathrm{B}\left(b^{\prime} \rightarrow Z b\right)=1$ |
| $>1230$ | 95 | 3 SIRUNYAN | 19BWCMS | $\mathrm{B}\left(b^{\prime} \rightarrow W t\right)=1$ |
| $>1350$ | 95 | 4 AABOUD 18 | 18AW ATLS | $\mathrm{B}\left(b^{\prime} \rightarrow W t\right)=1$ |
| $>1000$ | 95 | ${ }^{5}$ AABOUD 18 | 18CE ATLS | $\geq 2 \ell+E_{T}+\geq 1 b j$ |
| $>950$ | 95 | 6 AABOUD 1 | 18CL ATLS | $W t, Z b, h b$ modes |
| $>1010$ | 95 | 7,8 AABOUD 1 | 18CP ATLS | 2,3 $\ell$, singlet model |
| >1140 | 95 | 6,9 AABOUD 1 | 18CP ATLS | 2,3 , doublet model |
| >1220 | 95 | 10,11 AABOUD 18 | 18CR ATLS | singlet $b^{\prime}$. ATLAS Com nation |
| >1370 | 95 | 10,12 AABOUD 1 | 18CR ATLS | $b^{\prime}$ in a weak isospin do blet $\left(t^{\prime}, b^{\prime}\right)$. ATLAS combination. |
| $>910$ | 95 | 13 SIRUNYAN 1 | 18BM CMS | $W t, Z b, h b$ modes |
| > 845 | 95 | 14 SIRUNYAN 18 | 18Q CMS | $\mathrm{B}\left(b^{\prime} \rightarrow W u\right)=1$ |
| $>730$ | 95 | 15 SIRUNYAN 17 | 17AU CMS |  |
| $>880$ | 95 | 16 KHACHATRY... 1 | 16AN CMS | $\mathrm{B}\left(b^{\prime} \rightarrow W t\right)=1$ |
| $>620$ | 95 | 17 AAD 1 | 15BY ATLS | $W t, Z b, h b$ modes |
| $>730$ | 95 | 18 AAD 1 | 15BY ATLS | $\mathrm{B}\left(b^{\prime} \rightarrow W t\right)=1$ |
| $>810$ | 95 | 19 AAD 1 | 15z ATLS |  |
| $>755$ | 95 | 20 AAD 1 | 14AZ ATLS | $\mathrm{B}\left(b^{\prime} \rightarrow W t\right)=1$ |
| $>675$ | 95 | 21 CHATRCHYAN 1 | 13I CMS | $\mathrm{B}\left(b^{\prime} \rightarrow W t\right)=1$ |
| $>190$ | 95 | 22 ABAZOV 08x | 08x D0 | $\mathrm{c} \tau=200 \mathrm{~mm}$ |
| $>190$ | 95 | 23 ACOSTA 03 | 03 CDF | quasi-stable $b^{\prime}$ |

-     - We do not use the following data for averages, fits, limits, etc. - -

| <350, 580-635, > 700 | 95 |  | AAD | 15AR | ATLS | $\mathrm{B}\left(b^{\prime} \rightarrow H b\right)=1$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| > 690 | 95 |  | AAD | 15CN | ATLS | $\mathrm{B}\left(b^{\prime} \rightarrow W \mathrm{q}\right)=1(q=u)$ |
| $>480$ | 95 |  | AAD | 12AT | ATLS | $\mathrm{B}\left(b^{\prime} \rightarrow W t\right)=1$ |
| $>400$ | 95 | 27 | AAD | 12AU | ATLS | $\mathrm{B}\left(b^{\prime} \rightarrow Z b\right)=1$ |
| > 350 | 95 |  | AAD | 12BC | ATLS | $\begin{aligned} & \mathrm{B}\left(b^{\prime} \rightarrow W q\right)=1 \\ & \quad(q=u, c) \end{aligned}$ |
| $>450$ | 95 |  | AAD | 12BE | ATLS | $\mathrm{B}\left(b^{\prime} \rightarrow W t\right)=1$ |
| $>685$ | 95 |  | CHATRCHYAN | 12BH | CMS | $m_{t^{\prime}}=m_{b^{\prime}}$ |
| $>611$ | 95 |  | CHATRCHYAN | 12x | CMS | $\mathrm{B}\left(b^{\prime} \rightarrow W t\right)=1$ |
| > 372 | 95 |  | AALTONEN | 11J | CDF | $b^{\prime} \rightarrow W t$ |
| > 361 | 95 |  | CHATRCHYAN | 11L | CMS | Repl. by CHATRCHYAN $12 x$ |
| $>338$ | 95 |  | AALTONEN | 10 H | CDF | $b^{\prime} \rightarrow W$ t |
| > 380-430 | 95 | 35 | FLACCO | 10 | RVUE | $m_{b^{\prime}}>m_{t^{\prime}}$ |
| $>268$ | 95 | 36,37 | AALTONEN | 07C | CDF | $\mathrm{B}\left(b^{\prime} \rightarrow Z b\right)=1$ |
| > 199 | 95 |  | AFFOLDER | 00 | CDF | NC: $b^{\prime} \rightarrow Z b$ |


| $>148$ | 95 | 39 ABE | 98N | CDF | NC: $b^{\prime} \rightarrow Z \quad$ b + vertex |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $>96$ | 95 | 40 ABACHI | 97D | D0 | $\mathrm{NC}: b^{\prime} \rightarrow b \gamma$ |
| $>128$ | 95 | 41 ABACHI | 95F | D0 | $\ell \ell+$ jets, $\ell+$ jets |
| $>75$ | 95 | 42 MUKHOPAD.. | 93 | RVUE | NC: $b^{\prime} \rightarrow$ bll |
| $>85$ | 95 | 43 ABE | 92 | CDF | CC: $\ell \ell$ |
| $>72$ | 95 | 44 ABE | 90B | CDF | CC: $e+\mu$ |
| $>54$ | 95 | 45 AKESSON | 90 | UA2 | CC: $e+$ jets $+E_{T}$ |
| $>43$ | 95 | 46 ALBAJAR | 90B | UA1 | CC: $\mu+$ jets |
| $>34$ | 95 | 47 ALBAJAR | 88 | UA1 | CC: e or $\mu+$ jets |

${ }^{1}$ SIRUNYAN 20BI based on $137 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=13 \mathrm{TeV}$. Pair production of vector-like $b^{\prime}$ is seached for with each $b^{\prime}$ decaying into $Z b$ or $h b$. Analysis focuses on final states consisting of jets from six quarks. Mass limits are obtained for a variety of branching ratios of $b^{\prime}$ decays.
${ }^{2}$ SIRUNYAN 19AQ based on $35.9 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=13 \mathrm{TeV}$. Pair production of vector-like $b^{\prime}$ is seached for with one $b^{\prime}$ decaying into $Z b$ and the other $b^{\prime}$ decaying into $W t, Z b, h b$. Events with an opposite-sign lepton pair consistent with coming from $Z$ and jets are used. Mass limits are obtained for a variety of branching ratios of $b^{\prime}$.
${ }^{3}$ SIRUNYAN 19BW based on $35.9 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=13 \mathrm{TeV}$. The limit is for the pair-produced vector-like $b^{\prime}$ using all-hadronic final state. The analysis is made for the $Z b, W t, h b$ modes and mass limits are obtained for a variety of branching ratios.
${ }^{4}$ AABOUD 18AW based on $36.1 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=13 \mathrm{TeV}$. The limit is for the pair-produced vector-like $b^{\prime}$ using lepton-plus-jets final state. The search is also sensitive to the decays into $Z b$ and $H b$ final states.
${ }^{5}$ AABOUD 18CE based on $36.1 \mathrm{fb}^{-1}$ of proton-proton data taken at $\sqrt{s}=13 \mathrm{TeV}$. Events including a same-sign lepton pair are used. The limit is for a singlet model, assuming the branching ratios of $b^{\prime}$ into $Z b, W t$ and $H b$ as predicted by the model.
${ }^{6}$ AABOUD 18CL, AABOUD 18CP based on $36.1 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=13 \mathrm{TeV}$. The limit is for the pair-produced vector-like $b^{\prime}$ using all-hadronic final state. The analysis is particularly powerful for the $b^{\prime} \rightarrow h b$ mode. Assuming the pure decay only in this mode sets a limit $m_{b^{\prime}}>1010 \mathrm{GeV}$.
${ }^{7}$ AABOUD 18CP based on $36.1 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=13 \mathrm{TeV}$. Pair and single production of vector-like $b^{\prime}$ are seached for with at least one $b^{\prime}$ decaying into $Z b$. In the case of $\mathrm{B}\left(b^{\prime} \rightarrow Z b\right)=1$, the limit is $m_{b^{\prime}}>1220 \mathrm{GeV}$.
${ }^{8}$ The limit is for the singlet model, assuming that the branching ratios into $W t, Z b, h b$ add up to one.
${ }^{9}$ The limit is for the doublet model, assuming that the branching ratios into $W t, Z b, h b$ add up to one.
10 AABOUD 18CR based on $36.1 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=13 \mathrm{TeV}$. A combination of searches for the pair-produced vector-like $b^{\prime}$ in various decay channels $\left(b^{\prime} \rightarrow W t, Z b\right.$, $h b$ ). Also a model-independent limit is obtained as $m_{b^{\prime}}>1.03 \mathrm{TeV}$, assuming that the branching ratios into $Z b, W t$, and $h b$ add up to one.
11 The limit is for the singlet $b^{\prime}$.
12 The limit is for $b^{\prime}$ in a weak isospin doublet $\left(t^{\prime}, b^{\prime}\right)$ and $\left|V_{t^{\prime} b}\right| \ll\left|V_{t b^{\prime}}\right|$. For a $b^{\prime}$ in a doublet with a charge $-4 / 3$ vector-like quark, the limit $m_{b^{\prime}}>1.14 \mathrm{TeV}$ is obtained.
13 SIRUNYAN 18BM based on $35.9 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=13 \mathrm{TeV}$. The limit is for the pair-produced vector-like $b^{\prime}$. Three channels (single lepton, same-charge 2 leptons, or at least 3 leptons) are considered for various branching fraction combinations. Assuming $\mathrm{B}(t W)=1$, the limit is 1240 GeV and for $\mathrm{B}(b Z)=1$ it is 960 GeV .
14 SIRUNYAN 18Q based on $19.7 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=8 \mathrm{TeV}$. The limit is for the pair-produced vector-like $b^{\prime}$ that couple only to light quarks. Upper cross section limits
on the single production of a $b^{\prime}$ and constraints for other decay channels ( $Z q$ and $H q$ ) are also given in the paper.
15 SIRUNYAN 17AU based on $2.3-2.6 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=13 \mathrm{TeV}$. Limit on pairproduced singlet vector-like $b^{\prime}$ using one lepton and several jets. The mass bound is given for a $b^{\prime}$ transforming as a singlet under the electroweak symmetry group, assumed to decay through $W, Z$ or Higgs boson (which decays to jets) and to a third generation quark.
16 KHACHATRYAN 16AN based on $19.7 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=8 \mathrm{TeV}$. Limit on pairproduced vector-like $b^{\prime}$ using 1,2 , and $>2$ leptons as well as fully hadronic final states. Other limits depending on the branching fractions to $t W, b Z$, and $b H$ are given in Table IX.
${ }^{17}$ AAD 15BY based on $20.3 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=8 \mathrm{TeV}$. Limit on pair-produced vector-like $b^{\prime}$ assuming the branching fractions to $W, Z$, and $h$ modes of the singlet model. Used events containing $\geq 2 \ell+E_{T}+\geq 2 \mathrm{j}(\geq 1 b)$ and including a same-sign lepton pair.
18 AAD 15BY based on $20.3 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=8 \mathrm{TeV}$. Limit on pair-produced chiral $b^{\prime}$-quark. Used events containing $\geq 2 \ell+E_{T}+\geq 2 \mathrm{j}(\geq 1 b)$ and including a same-sign lepton pair.
${ }^{19}$ AAD $15 Z$ based on $20.3 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=8 \mathrm{TeV}$. Used events with $\ell+E_{T}+$ $\geq 6 \mathrm{j}(\geq 1 b)$ and at least one pair of jets from weak boson decay, primarily designed to select the signature $b^{\prime} \bar{b}^{\prime} \rightarrow W W t \bar{t} \rightarrow W W W W b \bar{b}$. This is a limit on pair-produced vector-like $b^{\prime}$. The lower mass limit is 640 GeV for a vector-like singlet $b^{\prime}$.
${ }^{20}$ Based on $20.3 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=8 \mathrm{TeV}$. No significant excess over SM expectation is found in the search for pair production or single production of $b^{\prime}$ in the events with dilepton from a high $p_{T} Z$ and additional jets ( $\geq 1 b$-tag). If instead of $\mathrm{B}\left(b^{\prime} \rightarrow W t\right)$ $=1$ an electroweak singlet with $\mathrm{B}\left(b^{\prime} \rightarrow W t\right) \sim 0.45$ is assumed, the limit reduces to 685 GeV .
21 Based on $5.0 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=7 \mathrm{TeV}$. CHATRCHYAN 13 looked for events with one isolated electron or muon, large $E_{T}$, and at least four jets with large transverse momenta, where one jet is likely to originate from the decay of a bottom quark.
${ }^{22}$ Result is based on $1.1 \mathrm{fb}^{-1}$ of data. No signal is found for the search of long-lived particles which decay into final states with two electrons or photons, and upper bound on the cross section times branching fraction is obtained for $2<\mathrm{c} \tau<7000 \mathrm{~mm}$; see Fig. 3. $95 \%$ CL excluded region of $b^{\prime}$ lifetime and mass is shown in Fig. 4.

23 ACOSTA 03 looked for long-lived fourth generation quarks in the data sample of 90 $\mathrm{pb}^{-1}$ of $\sqrt{s}=1.8 \mathrm{TeV} p \bar{p}$ collisions by using the muon-like penetration and anomalously high ionization energy loss signature. The corresponding lower mass bound for the charge $(2 / 3)$ e quark $\left(t^{\prime}\right)$ is 220 GeV . The $t^{\prime}$ bound is higher than the $b^{\prime}$ bound because $t^{\prime}$ is more likely to produce charged hadrons than $b^{\prime}$. The $95 \% \mathrm{CL}$ upper bounds for the production cross sections are given in their Fig. 3.
${ }^{24}$ AAD 15AR based on $20.3 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=8 \mathrm{TeV}$. Used lepton-plus-jets final state. See Fig. 24 for mass limits in the plane of $\mathrm{B}\left(b^{\prime} \rightarrow W t\right)$ vs. $\mathrm{B}\left(b^{\prime} \rightarrow H b\right)$ from $b^{\prime} \bar{b}^{\prime} \rightarrow H b+X$ searches.
${ }^{25}$ AAD 15CN based on $20.3 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=8 \mathrm{TeV}$. Limit on pair-production of chiral $b^{\prime}$-quark. Used events with $\ell+E_{T}+\geq 4 j$ (non- $b$-tagged). Limits on a heavy vector-like quark, which decays into $W q, Z q, h q$, are presented in the plane $\mathrm{B}(Q \rightarrow$ $W q)$ vs. $\mathrm{B}(Q \rightarrow h q)$ in Fig. 12.
${ }^{26}$ Based on $1.04 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=7 \mathrm{TeV}$. No signal is found for the search of heavy quark pair production that decay into $W$ and a $t$ quark in the events with a high $p_{T}$ isolated lepton, large $E_{T}$, and at least 6 jets in which one, two or more dijets are from $W$.
27 Based on $2.0 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=7 \mathrm{TeV}$. No $b^{\prime} \rightarrow Z b$ invariant mass peak is found in the search of heavy quark pair production that decay into $Z$ and a $b$ quark in
events with $Z \rightarrow e^{+} e^{-}$and at least one $b$-jet. The lower mass limit is 358 GeV for a vector-like singlet $b^{\prime}$ mixing solely with the third SM generation.
28 Based on $1.04 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=7 \mathrm{TeV}$. No signal is found for the search of heavy quark pair production that decay into $W$ and a quark in the events with dileptons, large $E_{T}$, and $\geq 2$ jets.
${ }^{29}$ Based on $1.04 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=7 \mathrm{TeV}$. AAD 12BE looked for events with two isolated like-sign leptons and at least 2 jets, large $E_{T}$ and $\mathrm{H}_{T}>350 \mathrm{GeV}$.
30 Based on $5 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=7 \mathrm{TeV}$. CHATRCHYAN 12BH searched for QCD and EW production of single and pair of degenerate 4 'th generation quarks that decay to $b W$ or $t W$. Absence of signal in events with one lepton, same-sign dileptons or trileptons gives the bound. With a mass difference of $25 \mathrm{GeV} / \mathrm{c}^{2}$ between $m_{t^{\prime}}$ and $m_{b^{\prime}}$, the corresponding limit shifts by about $\pm 20 \mathrm{GeV} / \mathrm{c}^{2}$.
31 Based on $4.9 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=7 \mathrm{TeV}$. CHATRCHYAN $12 \times$ looked for events with trileptons or same-sign dileptons and at least one $b$ jet.
32 Based on $4.8 \mathrm{fb}^{-1}$ of data in $p \bar{p}$ collisions at 1.96 TeV . AALTONEN 11 J looked for events with $\ell+E_{T}+\geq 5 j(\geq 1 b$ or $c)$. No signal is observed and the bound $\sigma\left(b^{\prime} \bar{b}^{\prime}\right)$ $<30 \mathrm{fb}$ for $m_{b^{\prime}}>375 \mathrm{GeV}$ is found for $\mathrm{B}\left(b^{\prime} \rightarrow W t\right)=1$.
33 Based on $34 \mathrm{pb}^{-1}$ of data in $p p$ collisions at 7 TeV . CHATRCHYAN 11L looked for multijet events with trileptons or same-sign dileptons. No excess above the SM background excludes $m_{b^{\prime}}$ between 255 and 361 GeV at $95 \% \mathrm{CL}$ for $\mathrm{B}\left(b^{\prime} \rightarrow W t\right)=1$.
${ }^{34}$ Based on $2.7 \mathrm{fb}^{-1}$ of data in $p \bar{p}$ collisions at $\sqrt{s}=1.96 \mathrm{TeV}$. AALTONEN 10 H looked for pair production of heavy quarks which decay into $t W^{-}$or $t W^{+}$, in events with same sign dileptons (e or $\mu$ ), several jets and large missing $E_{T}$. The result is obtained for $b^{\prime}$ which decays into $t W^{-}$. For the charge $5 / 3$ quark $\left(T_{5 / 3}\right)$ which decays into $t W^{+}$, $m_{T_{5 / 3}}>365 \mathrm{GeV}(95 \% \mathrm{CL})$ is found when it has the charge $-1 / 3$ partner B of the same mass.
${ }^{35}$ FLACCO 10 result is obtained from AALTONEN 10 H result of $m_{b^{\prime}}>338 \mathrm{GeV}$, by relaxing the condition $\mathrm{B}\left(b^{\prime} \rightarrow W t\right)=100 \%$ when $m_{b^{\prime}}>m_{t^{\prime}}$.
36 Result is based on $1.06 \mathrm{fb}^{-1}$ of data. No excess from the $\mathrm{SM} Z+$ jet events is found when $Z$ decays into ee or $\mu \mu$. The $m_{b^{\prime}}$ bound is found by comparing the resulting upper bound on $\sigma\left(b^{\prime} \bar{b}^{\prime}\right)$ [1-(1- $\left.\left.\mathrm{B}\left(b^{\prime} \rightarrow Z b\right)\right)^{2}\right]$ and the LO estimate of the $b^{\prime}$ pair production cross section shown in Fig. 38 of the article.
37 HUANG 08 reexamined the $b^{\prime}$ mass lower bound of 268 GeV obtained in AALTONEN 07C that assumes $\mathrm{B}\left(b^{\prime} \rightarrow Z b\right)=1$, which does not hold for $m_{b^{\prime}}>255 \mathrm{GeV}$. The lower mass bound is given in the plane of $\sin ^{2}\left(\theta_{t b^{\prime}}\right)$ and $m_{b^{\prime}}$.
38 AFFOLDER 00 looked for $b^{\prime}$ that decays in to $b+Z$. The signal searched for is $b b Z Z$ events where one $Z$ decays into $e^{+} e^{-}$or $\mu^{+} \mu^{-}$and the other $Z$ decays hadronically. The bound assumes $\mathrm{B}\left(b^{\prime} \rightarrow Z b\right)=100 \%$. Between 100 GeV and 199 GeV , the $95 \% \mathrm{CL}$ upper bound on $\sigma\left(b^{\prime} \rightarrow \bar{b}^{\prime}\right) \times \mathrm{B}^{2}\left(b^{\prime} \rightarrow Z b\right)$ is also given (see their Fig. 2).
39 ABE 98N looked for $Z \rightarrow e^{+} e^{-}$decays with displaced vertices. Quoted limit assumes $\mathrm{B}\left(b^{\prime} \rightarrow Z b\right)=1$ and $c \tau_{b^{\prime}}=1 \mathrm{~cm}$. The limit is lower than $m_{Z}+m_{b}(\sim 96 \mathrm{GeV})$ if $c \tau>22 \mathrm{~cm}$ or $c \tau<0.009 \mathrm{~cm}$. See their Fig. 4.
40 ABACHI 97D searched for $b^{\prime}$ that decays mainly via FCNC. They obtained $95 \%$ CL upper bounds on $\mathrm{B}\left(b^{\prime} \bar{b}^{\prime} \rightarrow \gamma+3\right.$ jets $)$ and $\mathrm{B}\left(b^{\prime} \bar{b}^{\prime} \rightarrow 2 \gamma+2\right.$ jets $)$, which can be interpreted as the lower mass bound $m_{b^{\prime}}>m_{Z}+m_{b}$.
${ }^{41} \mathrm{ABACHI} 95 \mathrm{~F}$ bound on the top-quark also applies to $b^{\prime}$ and $t^{\prime}$ quarks that decay predominantly into $W$. See FROGGATT 97.
${ }^{42}$ MUKHOPADHYAYA 93 analyze CDF dilepton data of ABE 92G in terms of a new quark decaying via flavor-changing neutral current. The above limit assumes $\mathrm{B}\left(b^{\prime} \rightarrow\right.$
$\left.b \ell^{+} \ell^{-}\right)=1 \%$. For an exotic quark decaying only via virtual $Z\left[\mathrm{~B}\left(b \ell^{+} \ell^{-}\right)=3 \%\right]$, the limit is 85 GeV .
43 ABE 92 dilepton analysis limit of $>85 \mathrm{GeV}$ at $\mathrm{CL}=95 \%$ also applies to $b^{\prime}$ quarks, as discussed in ABE 90B.
44 ABE 90B exclude the region $28-72 \mathrm{GeV}$.
${ }^{45}$ AKESSON 90 searched for events having an electron with $p_{T}>12 \mathrm{GeV}$, missing momentum $>15 \mathrm{GeV}$, and a jet with $E_{T}>10 \mathrm{GeV},|\eta|<2.2$, and excluded $m_{b^{\prime}}$ between 30 and 69 GeV .
${ }^{46}$ For the reduction of the limit due to non-charged-current decay modes, see Fig. 19 of ALBAJAR 90B.
47 ALBAJAR 88 study events at $E_{\mathrm{cm}}=546$ and 630 GeV with a muon or isolated electron, accompanied by one or more jets and find agreement with Monte Carlo predictions for the production of charm and bottom, without the need for a new quark. The lower mass limit is obtained by using a conservative estimate for the $b^{\prime} \bar{b}^{\prime}$ production cross section and by assuming that it cannot be produced in $W$ decays. The value quoted here is revised using the full $O\left(\alpha_{s}^{3}\right)$ cross section of ALTARELLI 88.

## $\boldsymbol{b}^{\prime}(-1 / 3)$ mass limits from single production in $p \bar{p}$ and $p p$ collisions

| $V A L U E(\mathrm{GeV})$ | CL\% | DOCUMENT ID | TECN | COMMENT |
| :---: | :---: | :---: | :---: | :---: |
| >1500 | 95 | 1 AAD 16AH | ATLS | $\begin{gathered} g b \rightarrow \underset{y}{ } \rightarrow{ }^{\prime} \rightarrow t W, \mathrm{~B}\left(b^{\prime} \rightarrow\right. \\ t W) \end{gathered}$ |
| >1390 | 95 | 2 KHACHATRY...16I | CMS | $\begin{gathered} g b \rightarrow \underset{ }{g} \rightarrow \\ t W)=1 \end{gathered}$ |
| >1430 | 95 | 3 KHACHATRY...16I | CMS | $\begin{gathered} g b \rightarrow \underset{y}{\prime} \rightarrow 1 \\ t W, \end{gathered}$ |
| >1530 | 95 | 4 KHACHATRY...16I | CMS | $\begin{gathered} g b \rightarrow \quad b^{\prime} \rightarrow t W, \mathrm{~B}\left(b^{\prime} \rightarrow\right. \\ t W)=1 \end{gathered}$ |
| $>693$ | 95 | ${ }^{5}$ ABAZOV 11F | D0 | $\begin{aligned} & q u \rightarrow q^{\prime} b^{\prime} \rightarrow q^{\prime}(W u) \\ & \quad \widetilde{\kappa}_{u b^{\prime}}=1, \mathrm{~B}\left(b^{\prime} \rightarrow W u\right)=1 \end{aligned}$ |
| $>430$ | 95 | 5 ABAZOV 11F | D0 | $\begin{aligned} & q d \rightarrow q b^{\prime} \rightarrow q(Z d) \\ & \quad \widetilde{\kappa}_{d b^{\prime}}=\sqrt{2}, \mathrm{~B}\left(b^{\prime} \rightarrow Z d\right)=1 \end{aligned}$ |

-     - We do not use the following data for averages, fits, limits, etc. - - -

6 SIRUNYAN 19AI CMS $b Z / t W \rightarrow b^{\prime} \rightarrow t W$
${ }^{1}$ AAD 16 AH based on $20.3 \mathrm{fb}^{-1}$ of data in $p p$ collisions at 8 TeV . No significant excess over SM expectation is found in the search for a vector-like $b^{\prime}$ in the single-lepton and dilepton channels $(\ell$ or $\ell \ell)+1,2,3 j(\geq 1 b)$. The model assumes that the $b^{\prime}$ has the excited quark couplings.
${ }^{2}$ Based on $19.7 \mathrm{fb}^{-1}$ of data in $p p$ collisions at 8 TeV . Limit on left-handed $b^{\prime}$ assuming $100 \%$ decay to $t W$ and using all-hadronic, lepton + jets, and dilepton final states.
${ }^{3}$ Based on $19.7 \mathrm{fb}^{-1}$ of data in $p p$ collisions at 8 TeV . Limit on right-handed $b^{\prime}$ assuming $100 \%$ decay to $t W$ and using all-hadronic, lepton + jets, and dilepton final states.
${ }^{4}$ Based on $19.7 \mathrm{fb}^{-1}$ of data in $p p$ collisions at 8 TeV . Limit on vector-like $b^{\prime}$ assuming $100 \%$ decay to $t W$ and using all-hadronic, lepton + jets, and dilepton final states.
${ }^{5}$ Based on $5.4 \mathrm{fb}^{-1}$ of data in ppbar collisions at 1.96 TeV . ABAZOV 11 F looked for single production of $b^{\prime}$ via the $W$ or $Z$ coupling to the first generation up or down quarks, respectively. Model independent cross section limits for the single production processes $p \bar{p} \rightarrow b^{\prime} q \rightarrow W u q$, and $p \bar{p} \rightarrow b^{\prime} q \rightarrow Z d q$ are given in Figs. 3 and 4, respectively, and the mass limits are obtained for the model of ATRE 09 with degenerate bi-doublets of vector-like quarks.
${ }^{6}$ SIRUNYAN 19AI based on $35.9 \mathrm{fb}^{-1}$ of $p p$ data at $\sqrt{s}=13 \mathrm{TeV}$. Exclusion limits are set on the product of the production cross section and branching fraction for the $b^{\prime}(-1 / 3)+b$ and $b^{\prime}(-1 / 3)+t$ modes as a function of the vector-like quark mass in Figs. 7 and 8 and in Tab. 2 for relative vector-like quark widths between 1 and $30 \%$ for
left- and right-handed vector-like quark couplings. No significant deviation from the SM prediction is observed

## MASS LIMITS for $\boldsymbol{b}^{\prime}$ (4 $4^{\text {th }}$ Generation) Quark or Hadron in $\boldsymbol{e}^{+} \boldsymbol{e}^{-}$Collisions

Search for hadrons containing a fourth-generation $-1 / 3$ quark denoted $b^{\prime}$.
The last column specifies the assumption for the decay mode (C C denotes the conventional charged-current decay) and the event signature which is looked for.

| VALUE (GeV) | CL\% | DOCUMENT ID |  | TECN | $\frac{\text { COMMENT }}{\text { any decay }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| >46.0 | 95 | 1 DECAMP | 90F | ALEP |  |
| - - We do not use the following data for averages, fits, limits, etc. - - |  |  |  |  |  |
| none 96-103 | 95 | ${ }^{2}$ ABDALLAH <br> 3 ADRIANI | $\begin{aligned} & 07 \\ & 93 \mathrm{G} \end{aligned}$ | $\begin{aligned} & \text { DLPH } \\ & \text { L3 } \end{aligned}$ | $b^{\prime} \rightarrow b Z, c W$ <br> Quarkonium |
| $>44.7$ | 95 | ADRIANI | 93M | L3 | $\Gamma(Z)$ |
| $>45$ | 95 | ABREU | 91F | DLPH | $\Gamma(Z)$ |
| none 19.4-28.2 | 95 | ABE | 90D | VNS | Any decay; event shape |
| $>45.0$ | 95 | ABREU | 90D | DLPH | $B(C C)=1$; event shape |
| $>44.5$ | 95 | ${ }^{4}$ ABREU | 90D | DLPH | $\begin{aligned} b^{\prime} & \rightarrow c H^{-}, H^{-} \rightarrow \\ & \bar{c} s, \tau^{-} \nu \end{aligned}$ |
| $>40.5$ | 95 | ${ }^{5}$ ABREU | 90D | DLPH | $\Gamma(Z \rightarrow$ hadrons $)$ |
| $>28.3$ | 95 | ADACHI | 90 | TOPZ | $\begin{aligned} & \mathrm{B}(\mathrm{FCNC})=100 \% \text {; isol. } \\ & \quad \gamma \text { or } 4 \text { jets } \end{aligned}$ |
| $>41.4$ | 95 | 6 AKRAWY | 90B | OPAL | Any decay; acoplanarity |
| $>45.2$ | 95 | 6 AKRAWY | 90B | OPAL | $\mathrm{B}(C C)=1 ; \text { acopla- }$ narity |
| $>46$ | 95 | ${ }^{7}$ AKRAWY | 90J | OPAL | $b^{\prime} \rightarrow \gamma+$ any |
| $>27.5$ | 95 | ${ }^{8} \mathrm{ABE}$ | 89E | VNS | $\mathrm{B}(C C)=1 ; \mu, e$ |
| none 11.4-27.3 | 95 | 9 ABE | 89G | VNS | $\begin{aligned} & \mathrm{B}\left(b^{\prime} \rightarrow b \gamma\right)>10 \% \\ & \quad \text { isolated } \gamma \end{aligned}$ |
| $>44.7$ | 95 | 10 ABRAMS | 89C | MRK2 | $B(C C)=100 \%$; isol. track |
| $>42.7$ | 95 | 10 ABRAMS | 89C | MRK2 | $\begin{aligned} & \mathrm{B}(b g)=100 \% ; \text { event } \\ & \text { shape } \end{aligned}$ |
| $>42.0$ | 95 | 10 ABRAMS | 89C | MRK2 | Any decay; event shape |
| $>28.4$ | 95 | 11,12 ADACHI | 89C | TOPZ | $\mathrm{B}(\mathrm{C} C)=1 ; \mu$ |
| $>28.8$ | 95 | 13 ENO | 89 | AMY | $\mathrm{B}(\mathrm{CC}) \geq 90 \% ; \mu, e$ |
| $>27.2$ | 95 | 13,14 ENO | 89 | AMY | any decay; event shape |
| $>29.0$ | 95 | 13 ENO | 89 | AMY |  |
| $>24.4$ | 95 | 15 IGARASHI | 88 | AMY | $\mu, e$ |
| $>23.8$ | 95 | 16 SAGAWA | 88 | AMY | event shape |
| $>22.7$ | 95 | 17 ADEVA | 86 | MRKJ | $\mu$ |
| $>21$ |  | 18 ALTHOFF | 84C | TASS | $R$, event shape |
| >19 |  | 19 ALTHOFF | 841 | TASS | Aplanarity |

${ }^{1}$ DECAMP 90F looked for isolated charged particles, for isolated photons, and for four-jet final states. The modes $b^{\prime} \rightarrow b g$ for $\mathrm{B}\left(b^{\prime} \rightarrow b g\right)>65 \% b^{\prime} \rightarrow b \gamma$ for $\mathrm{B}\left(b^{\prime} \rightarrow b \gamma\right)$ $>5 \%$ are excluded. Charged Higgs decay were not discussed.
${ }^{2}$ ABDALLAH 07 searched for $b^{\prime}$ pair production at $E_{\mathrm{cm}}=196-209 \mathrm{GeV}$, with $420 \mathrm{pb}^{-1}$. No signal leads to the $95 \% \mathrm{CL}$ upper limits on $\mathrm{B}\left(b^{\prime} \rightarrow b Z\right)$ and $\mathrm{B}\left(b^{\prime} \rightarrow c W\right)$ for $m_{b^{\prime}}$ $=96$ to 103 GeV .
${ }^{3}$ ADRIANI 93 search for vector quarkonium states near $Z$ and give limit on quarkonium$Z$ mixing parameter $\delta m^{2}<(10-30) \mathrm{GeV}^{2}(95 \% \mathrm{CL})$ for the mass $88-94.5 \mathrm{GeV}$. Using Richardson potential, a $1 \mathrm{~S}\left(b^{\prime} \bar{b}^{\prime}\right)$ state is excluded for the mass range $87.7-94.7 \mathrm{GeV}$. This range depends on the potential choice.
${ }^{4}$ ABREU 90D assumed $m_{H^{-}}<m_{b^{\prime}}-3 \mathrm{GeV}$.
${ }^{5}$ Superseded by ABREU 91F.
${ }^{6}$ AKRAWY 90 B search was restricted to data near the $Z$ peak at $E_{\mathrm{cm}}=91.26 \mathrm{GeV}$ at LEP. The excluded region is between 23.6 and 41.4 GeV if no $\mathrm{H}^{+}$decays exist. For charged Higgs decays the excluded regions are between ( $m_{H^{+}}+1.5 \mathrm{GeV}$ ) and 45.5 GeV .
${ }^{7}$ AKRAWY 90」 search for isolated photons in hadronic $Z$ decay and derive $\mathrm{B}\left(Z \rightarrow b^{\prime} \bar{b}^{\prime}\right) \cdot \mathrm{B}\left(b^{\prime} \rightarrow \gamma \mathrm{X}\right) / \mathrm{B}(Z \rightarrow$ hadrons $)<2.2 \times 10^{-3}$. Mass limit assumes $\mathrm{B}\left(b^{\prime} \rightarrow \gamma \mathrm{X}\right)>10 \%$.
${ }^{8}$ ABE 89E search at $E_{\mathrm{cm}}=56-57 \mathrm{GeV}$ at TRISTAN for multihadron events with a spherical shape (using thrust and acoplanarity) or containing isolated leptons.
${ }^{9}$ ABE 89G search was at $E_{\mathrm{cm}}=55-60.8 \mathrm{GeV}$ at TRISTAN.
${ }^{10}$ If the photonic decay mode is large $\left(\mathrm{B}\left(b^{\prime} \rightarrow b \gamma\right)>25 \%\right)$, the ABRAMS 89C limit is 45.4 GeV . The limit for for Higgs decay ( $b^{\prime} \rightarrow c \mathrm{H}^{-}, \mathrm{H}^{-} \rightarrow \bar{c} s$ ) is 45.2 GeV .
${ }^{11}$ ADACHI 89C search was at $E_{\mathrm{cm}}=56.5-60.8 \mathrm{GeV}$ at TRISTAN using multi-hadron events accompanying muons.
${ }^{12}$ ADACHI 89C also gives limits for any mixture of $C C$ and $b g$ decays.
${ }^{13}$ ENO 89 search at $E_{\mathrm{cm}}=50-60.8$ at TRISTAN.
${ }^{14}$ ENO 89 considers arbitrary mixture of the charged current, $b g$, and $b \gamma$ decays.
${ }^{15}$ IGARASHI 88 searches for leptons in low-thrust events and gives $\Delta R\left(b^{\prime}\right)<0.26$ (95\% CL ) assuming charged current decay, which translates to $m_{b^{\prime}}>24.4 \mathrm{GeV}$.
${ }^{16}$ SAGAWA 88 set limit $\sigma$ (top $)<6.1 \mathrm{pb}$ at $\mathrm{CL}=95 \%$ for top-flavored hadron production from event shape analyses at $E_{\mathrm{cm}}=52 \mathrm{GeV}$. By using the quark parton model crosssection formula near threshold, the above limit leads to lower mass bounds of 23.8 GeV for charge $-1 / 3$ quarks.
17 ADEVA 86 give $95 \% \mathrm{CL}$ upper bound on an excess of the normalized cross section, $\Delta R$, as a function of the minimum c.m. energy (see their figure 3). Production of a pair of $1 / 3$ charge quarks is excluded up to $E_{\mathrm{cm}}=45.4 \mathrm{GeV}$.
${ }^{18}$ ALTHOFF 84 C narrow state search sets limit $\Gamma\left(e^{+} e^{-}\right) \mathrm{B}$ (hadrons) $<2.4 \mathrm{keV} \mathrm{CL}=95 \%$ and heavy charge $1 / 3$ quark pair production $m>21 \mathrm{GeV}, \mathrm{CL}=95 \%$.
${ }^{19}$ ALTHOFF 84 I exclude heavy quark pair production for $7<m<19 \mathrm{GeV}$ ( $1 / 3$ charge) using aplanarity distributions ( $\mathrm{CL}=95 \%$ ).

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